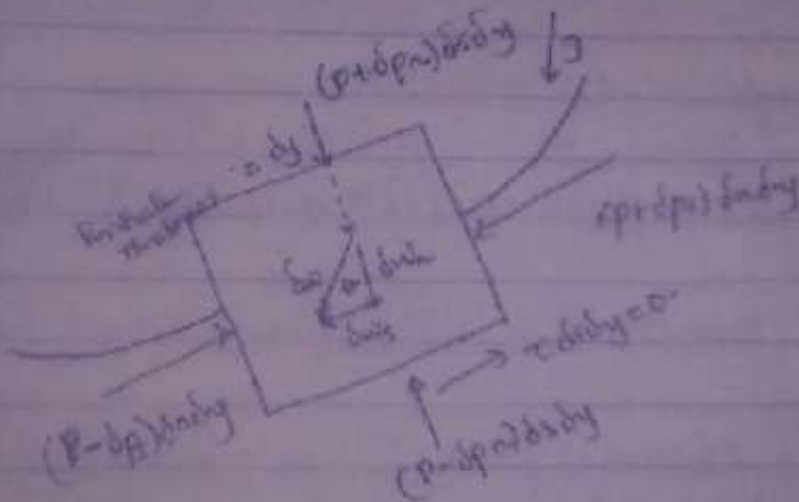


Component of weight force in the direction of stream line direction,  $\delta z$

$$\sum \delta F_z = \delta m g_z = \delta m g \frac{\delta z}{\delta l} = \rho \delta V g \frac{\delta z}{\delta l}$$

Component of weight force in the direction of stream line  $\delta l$

$$\delta F_z = -\delta m g \cos \theta = -\gamma \delta V \cos \theta$$



Use the 1st order Taylor series expansion for the pressure field:

$$\Delta p_s \approx \frac{\partial p}{\partial s} \cdot \frac{\delta s}{2}$$

The net pressure force on the particle in the stream line direction

$$\begin{aligned} \delta F_p &= (p - \Delta p) \delta n \delta y - (p + \Delta p) \delta n \delta y = -2 \Delta p_s \delta n \delta y \\ &= -\frac{\partial p}{\partial s} \delta s \delta n \delta y = +\frac{\partial p}{\partial s} \delta V \end{aligned}$$

net force acting in the streamline direction on the whole is

$$\begin{aligned}\sum \delta F_s &= \delta W_s + \delta F_{ps} = \left[ -\gamma \sin \theta - \frac{dp}{ds} \right] \delta V \\ &= -\gamma \sin \theta - \frac{dp}{ds} = \rho V \frac{dV}{ds} = \rho a_s\end{aligned}$$

Noting that  $\sin \theta = \frac{dz}{ds}$  and  $V \frac{dV}{ds} = \frac{1}{2} \frac{dV^2}{ds}$ , the

above eqn can be rearranged and integrated:

$$-\gamma \frac{dz}{ds} - \frac{dp}{ds} = \frac{1}{2} \rho \frac{dV^2}{ds}$$

$$\delta p + \frac{1}{2} \rho d(V^2) + \gamma \delta z = 0$$

(along a stream  
line)

For constant acceleration of gravity

$$\int \frac{dp}{\rho} + \frac{1}{2} V^2 + gz = C$$

(along a stream  
line)

For steady, inviscid and incompressible <sup>flow</sup> fluid, we have the celebrated Bernoulli's Equation

$$\boxed{P + \frac{1}{2} \rho V^2 + \gamma z = C} \quad (\text{along stream line})$$

$$P_1 + \frac{1}{2} \rho_1 V_1^2 + \gamma_1 z_1 = P_2 + \frac{1}{2} \rho_2 V_2^2 + \gamma_2 z_2$$